Damping and complex modes

- 1. Devices
- 2. Coupling/energy levels
- 3. Damping mechanisms, constitutive laws
- 4. Reduced models for damping: from material to structure





MS2SC PROVIR http://savoir.ensam.eu/moodle/course/view.php?id=1874 http://savoir.ensam.eu/moodle/course/view.php?id=490



Tuned mass dampers / vibration absorber





I wo lower resonances
One anti-resonance



- Countless applications : helicopters, buildings, lamp posts, cars, ...
- Current trend : non-linear absorber (self tuning)

Sample dampers

- Fluid sloshing : No moving parts
- Rings, cables, side masses, ...
- Electrical wires













www.multitech-fr.com

Damping due to local modes







Soize "structural fuzzy" 1987 Arnoux, A. Batou, C. Soize, L. Gagliardini. JSV 2013 Loukota, Passieux, Michon. Journal of Aircraft 2017

Damping devices

- Resonant mass
- Struts, point connections
- Line/surface treatments

Q2 : undamped nominal + coupling



Energy coupling & NL system



Damping occurs in the transition





FEMTo-ST Orion Testbed





- Contact with broadband excitation level
- Friction : amplitude/normal force dependence
- Viscoelasticity : temperature/frequency dependence



ENSAM/PhD. Hammami 2014 Hammami/Balmes/Guskov, MSSP 2015

A mechanical root locus (coupling)



Related ideas : electro-mechanical coupling coefficient, pole/zero distance, modal strain energy

Coupling & mode crossing

- Variable contact surface, contact, sliding can generate coupling
- Horizontal lines : global modes
- S curve : local cable guide mode





Sensors can help But can we track modes better ?

Lab work 3 : tracking & reanalysis

Q3 : Damping mechanisms

Damping : mechanism by which the energy in the considered system is diminished

- Non elastic behavior
 - Material : viscoelastic, plastic, ...
 - Joint friction
 - Inter-frequency coupling through non-linearities
- Coupling with other systems
 - Fluid/structure, particle, soil, thermal, ... interactions
 - Active control systems (piezo, electromagnetic, ...)
 - Rotating parts

Linear, Non-linear

Viscoelasticity = LTI material

- Viscoelasticity (Bolzmannian material), stress f(strain history) linear in amplitude
- Viscous relaxation : convolution with displacement or velocity (non-parametric)

$$F = \int_{-\infty}^{+\infty} K(t-r)x(r)dr = \int_{-\infty}^{+\infty} R(t-r)\dot{x}(r)dr$$

- Parametric model representation (Maxwell, rational fraction)

$$K(s) = K^{0} \prod_{i=1}^{N} \frac{1 - \frac{s}{z^{i}}}{1 - \frac{s}{p^{i}}} = K^{\infty} \left(1 + \sum_{i=1}^{N} -\frac{g^{i}}{\frac{s}{\omega^{i}}} \right) = K^{\infty} \left(g^{0} + \sum_{i=1}^{N} -\frac{g^{i}s}{s + \omega^{i}} \right)$$

- Differential equation formulation (stress or stress rate relaxation)

$$\frac{\dot{F}^i}{\omega^i} + F^i = -g^i F^\infty(x)$$
, or $\dot{F}^i + \omega^i F^i = g^i \dot{F}^\infty(x)$

- Order selection: place poles and zeros in frequency domain (> 1/decade)



F. Renaud & al. : A new identification method of viscoelastic behavior DOI 10.1016/j.ymssp.2010.09.002 Rafael Penas, PhD ENSAM 2021, MSSP 2022

Viscoelasticity 2 : frequency domain

- Time convolution $\sigma(t) = \int_{-\infty}^{t} C(t-\tau)\epsilon(\tau)$
- Frequency complex modulus $\sigma = E(\omega, T, \epsilon_S)\epsilon$
- $E = E'(1 + i\eta)$
- *E'* storage modulus, $\eta = \tan(\delta)$ loss factor
- Temperature/frequency equivalence

 $E(\omega, T, \epsilon_S) = E(\omega\alpha(T, \epsilon_S))$ generalized coordinate in Wicket plot



[3] SDT/Visc www.sdtools.com/pdf/visc.pdf

Temperature

10

SSO 0.5



Modulus / loss map of common materials



 $\sigma(\omega) = E_{complex} \epsilon(\omega) = E_{real} (1 + i\eta(\omega))\epsilon = E_{real} (1 + i \tan(\delta)(\omega))\epsilon$

Hysteretic damping : keywords

- Stress = fct (strain history) but independent on rate
- Non parametric : triangular testing
 - Masing rule (homotethy + offset)
 - Madelung assumption (loop closure)
 - Payne effect (E \downarrow amplitude \uparrow)
 - Mullins (first loading adaptation)
- Parametric models : Iwan, Dahl, Berg

[1] D. J. Segalman, "A Four-Parameter Iwan Model for Lap-Type Joints," Journal of Applied Mechanics, vol. 72, no. 5, pp. 752-760, Sep. 2005
 [2] B. Bourgeteau, "Modélisation numérique des articulations en caoutchouc de la liaison au sol automobile en simulation multi-corps transitoire," Ecole Centrale Paris, 2009

[3] R. Penas, " Models for dissipative bushings in multibody dynamics ", ENSAM, 2021

 $\times 10^4$ σ_{1,2} [Pa] -0.05 0 0.05 e [1] -0.6 -0.4 -0.2 0 02 04 €[1]



Rate independent hysteresis

- Hysteretic relaxation (non-parametric) $F(x) = F(x_{turn}) + \int_{0}^{|x-x_{turn}|} K_{f}(r) dr$
- Selected order parametric representation Jenkins cells, in an Iwan series arrangement
- Force formulation

$$F^i = \frac{g^i}{g^0} F^0$$
, if $F^i \operatorname{sign}(\dot{x}) < F^i_f$, and $F^i = \operatorname{sign}(\dot{x}) F^i_f$, if $\frac{g^i}{g^0} |F^0| \ge F^i_f$

Displacement formulation

$$\dot{x}^i = 0$$
, if $|x - x^i| < rac{F_f^i}{K^i}$, and $\dot{x}^i = \dot{x}$, if $|x - x^i| < rac{F_f^i}{K^i}$

- Definition of break points and stiffness

$$K_b^k = \sum_{i < k} K^i$$
, and $x_b^k = x_{turn} + \frac{F_f^k}{K^k}$

- Order selection: place x_b and K_b in the hysteretic relaxation curve





Parallel between viscoelasticity & hysteresis

- Fit non-parametric curves : discretize with selected order







- Identify order independent parameters

• STS
$$F - F_{turn} = \left(\sqrt{NC|x - x_{turn}| + \frac{NC}{2\hat{K}^0}} - \frac{NC}{2\hat{K}^1}\right) \operatorname{sign}(x - x_{turn})$$

- Segalman (4 parameter friction)
- Fractional derivative $F(s) + \frac{1}{\omega_c} \frac{\mathrm{d}^{\alpha} F}{\mathrm{d} t^{\alpha}} = K^0 x + \frac{1}{\omega_c^{\alpha}} K^{\infty} \frac{\mathrm{d}^{\alpha} x}{\mathrm{d} t^{\alpha}}$



• Implementation requires discretization = order selection

Adding hyperelasticity & combination

Material/mount model split in static, path and dynamic effects

• Non-parametric models for: hyperelasticity, hysteresis and viscoelasticity $E(\omega, A, \epsilon_0)$







$$\dot{F}^{i} + \left(\omega^{i} + \frac{|\dot{x}|}{x_{f}^{i}}\right)F^{i} = \frac{g^{i}}{g_{0}}\dot{F}^{0}(x)$$

- Among all cell models that combine HE, HYST and VE
 - stress rate relaxation $\dot{F}^0(x)$ is needed
 - NL relaxation times $\omega^{i} + \frac{|\dot{x}|}{x_{c}^{i}}$ explain Paine effect

20

Contact/friction

12

Normalised E₀ [N/mm³]

0

1

2

3 4

5

Contact Pressure [MPa]

6 7 8 9

10

- Surface contact/friction model
- Idealization Signorini/Coulomb
- $\begin{cases} \begin{bmatrix} u_n \end{bmatrix} \le 0 \\ R_n \le 0 \\ R_n \begin{bmatrix} u_n \end{bmatrix} = 0 \end{cases} \qquad R_t = -\mu |R_n| \frac{\begin{bmatrix} \dot{u}_t \end{bmatrix}}{\begin{bmatrix} \dot{u}_t \end{bmatrix}}$

- Reality
 - micro-scale effects
 - structural effects
 - $F_N(gap)$ and F_T hysteretic + dependent on F_N







Q4 : from material to structure model



- [1] Hammami, Balmes, Guskov, « Numerical design and test on an assembled structure of a bolted joint with viscoelastic damping », *MSSP*, v70, p.714–724, 2015
- [2] Farhat, Avery, Chapman, Cortial, « Dimensional reduction of nonlinear finite element dynamic models with finite rotations and energy-based mesh sampling and weighting for computational efficiency», *IJNME*, vol. 98, nº 9, p. 625-662, 2014.
- [3] F. Casenave, N. Akkari, F. Bordeu, C. Rey, and D. Ryckelynck, "A nonintrusive distributed reduced-order modeling framework for nonlinear structural mechanics—Application to elastoviscoplastic computations," IJNME, vol. 121, no. 1, pp. 32–53, Jan. 2020, doi: 10.1002/nme.6187.

Historical method : V_0 = reduction modes (poor)

MSE (modal strain energy) is a kinematic reduction using elastic modes and general damping

$$\{q\} = [T] \{q_r\} = [\phi_1 ... \phi_{NM}] \{q_r\}$$
$$\left[s^2 [I] + [\phi]^T [Im(Z(s))] [\phi] + \left[{}^{\backslash} \omega_{j_{\backslash}}^2 \right] \right] \{p\} (s) = [\phi]^T \{F(s)\}$$

NASTRAN

- SOL110 modal complex modes
- SOL111 modal frequency response
- SOL112 modal transient

When does MSE work?



"Real modes"	Residues for I/O pairs line up
"Complex modes"	Residues have a phase spread
Poor modes	Have complex residues

When are complex modes nearly real?

• Uncoupling criterion (Hasselman) \Leftrightarrow

$$\min(\zeta_1\omega_1,\zeta_2\omega_2)/|\omega_1-\omega_2|\ll 1$$

corresponds to non overlap of peaks : using real modes is then OK

Proof based on damping matrix positiveness

 $\gamma_{12}\gamma_{21}/(\gamma_{11}\gamma_{22}) < 1$

When damping is enhanced by design, modes are often complex

Historical V_1 = reduction modes + enrichment (good)



- •Modes (MSE)
- Multi-model
- Modes + static

 $T=[\Phi(p_0)]$ $T=[\Phi(p_1) \Phi(p_2)]$ $T=[\Phi(p_0) K_o^{-1} [Im(Z-Zo)] \Phi(p_0)]$

Reduction Bases/Starting Point

- Modal T=[$\Phi(p_0)$]
- Modal + static Responses to input $T=[\Phi(p_0) K^{-1}b]$
- Modal + static responses to representative loads $T=[\Phi(p_0) K^{-1}[b R_j] \Phi(p_k)]$
- Frequencies of corrections
 differ from modes
 : the added information is VERY different



MSE validity FRF prediction



Eigenvalue extraction

- Fixed E sparse solvers are not robust and slow \Rightarrow reduction T + full solver $[Z(E_i, \lambda_j)]\{\psi_j\} = \{0\}$
- Analytic E(s) not common + associated sparse eigenvalue not available ⇒ only used for transient
- Need to track $\{\psi_j\}, \omega_j, \zeta_j = f(E(T, \omega_j), \eta(T, \omega_j))$ \Rightarrow identify from forced response \Rightarrow interpolate from $f(E, \eta)$



Structural mechanisms

- Compression & shear springs
 - Integration gives orders of magnitude choice in functional stiffness
 - Possibly structured junctions
- Free layer bending (compression)
 - Inefficient if damping material is too soft
- Constrained layer bending (shear)
- Trough thickness dynamics
 - Only efficient in resonant (high frequency behavior)









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Optimization is necessary

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- Design space : material, thickness, surface (GS/h or EA/h)
- Multiple orders of magnitude Damping ratio target example



Why optimization is a challenge

- Large models (x1e6 DOF)
- Many modes /frequencies to compute (x1e3)
- Multiple designs (x10)
- Large parametric variations (1e4)
- Complex objectives (x1e3 disp)





Sum(Im(H))

Historical V_2 = multi-model & CMS

Objectives:

- complex mode stability
- Transient simulation







Historical V_3 = hyper reduction of transient NL

HR phase 1 : generic form of FEM problems

- 1. FEM kinematics $u(x,t) = [N(x)]\{q(t)\}$
- 2. Material evolution : $S = f_{mat}(\nabla u, u_{int})$
- 3. Virtual work $M\ddot{q} + \mathcal{F}_{int}(q, u_{int}) = \mathcal{F}_{ext}$
- Integration at quadrature points g (disassembly)

$$\mathcal{F}_{int}(q, u_{int}) = \sum_{g} \mathbb{C}^{T} J_{g} w_{g} f_{mat}(\epsilon, u_{int})$$
$$= [\mathbb{B}]_{N \times (N_{g} \times N_{f})} \left\{ f_{mat} \left([\mathbb{C}]_{(N \epsilon \times N_{g}) \times N} \{q\}, u_{int} \right) \right\}$$



HR phase 2 : kinematic reduction

- 1/ Off-line learning (non-linear High Fidelity Simulation) Snapshots or iterative method
- 2/ Kinematic reduction using SVD, or CMS





Reduced equations of motion $T^T M T \ddot{q}_r + [T^T \mathbb{B}]_{Nqr \times Ng} f_{mat}(\mathbb{C}T q_r, u_{int}) = T^T F_{ext}$

- N_g (gauss points) remains large
- $f_{mat}(\mathbb{C}Tq_r, u_{int})$ evaluation dominates

HR phase 3 : operator reduction

3/ Operator reduction (hyper-red.) = less points + new weights $\{F_{int}\} = \int_{\Omega} f_{mat}(x_g, t) \, \mathrm{d}V \approx \sum_{g} \mathbb{C}^{T} J_{g} w_{g} f_{mat g} = \mathbb{B}(x_g) f_{mat g}(q, t)$

Choose NT learning points and minimize $\left\|w_g^*\right\|_0$ subject to^[1]

 $\left\| \left[\left[T^T \mathbb{B} \right] S_l \right]_{(NT \times NR)} - \left[\left[T^T \mathbb{B}_g \right] S_{lg} \right]_{(NT \times NR) \times Ng} \{ w_g^* \} \right\|_2 < \varepsilon_{tol} \text{ and } w_g^* > 0$





- [2] F. Casenave, N. Akkari, F. Bordeu, C. Rey, and D. Ryckelynck, "A nonintrusive distributed reduced-order modeling framework for nonlinear structural mechanics—Application to elastoviscoplastic computations," IJNME, vol. 121, no. 1, pp. 32–53, Jan. 2020, doi: 10.1002/nme.6187
- [3] Penas, Rafael, Models of dissipative bushing in multibody dynamics, PhD ENSAM 2021





System = macro / micro

Field distribution : macroscopic law not a simple function • of material behavior



Run Number : **